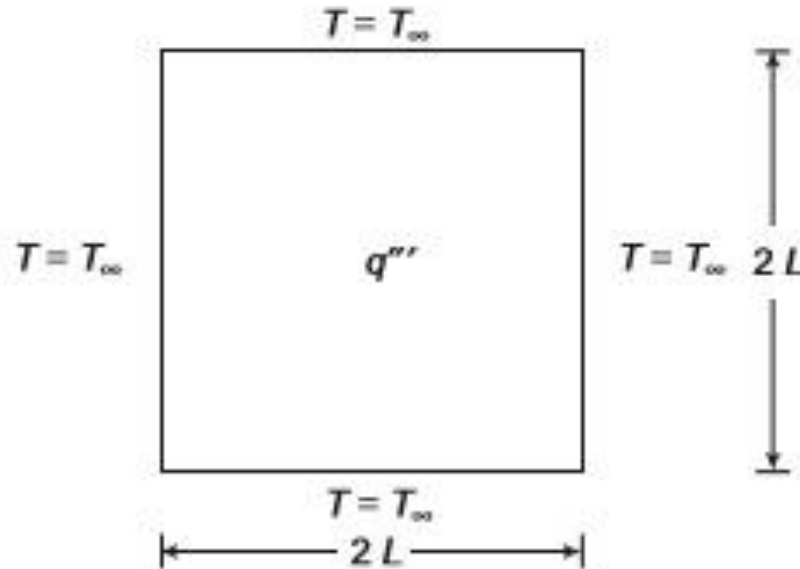


# ELEMENTS OF COMPUTATIONAL FLUID DYNAMICS

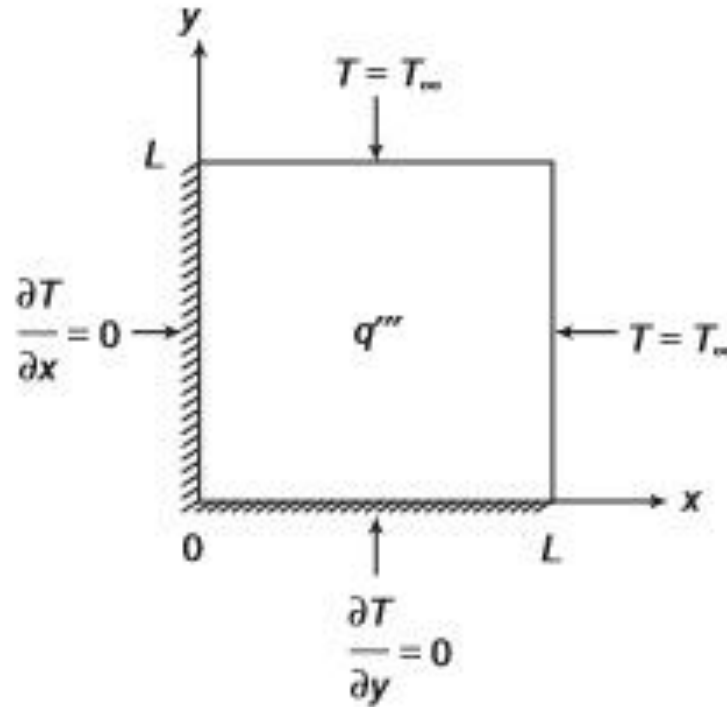
## Chapter - 7

## Two-dimensional steady-state problem



- ✓ The case of steady heat conduction in a long square slab ( $2L \times 2L$ ) in which heat is generated at a uniform rate of  $q'''$  W/m<sup>3</sup>.
- ✓ The problem can be assumed to be a two dimensional as the dimension of the slab is much longer in the direction normal to the cross-sectional plane; therefore, end effects can be neglected.
- ✓ All four sides are maintained at  $T = T_{\infty}$ , temperature of the surrounding fluid, assuming a large heat transfer coefficient.

## Consideration of symmetry



- A close look at the physics of the problem reveals that the problem is geometrically and thermally symmetric.
- Therefore, from the temperature distribution in any quarter of the physical domain by mirror-imaging, one can get the solution for the entire region.
- The use of symmetry enables the numerical analyst to obtain the solution much faster as the number of grid points is greatly reduced.

## Governing differential equation

The energy equation at the steady state (assuming constant  $k$ )

$$k \left( \frac{\partial^2 T}{\partial x^2} + \frac{\partial^2 T}{\partial y^2} \right) + q''' = 0$$

## Boundary conditions

Boundary conditions are

$$BC1: \text{ at } x = 0, \quad \frac{\partial T}{\partial x} = 0$$

$$BC2: \text{ at } x = L, \quad T = T_\infty$$

$$BC3: \text{ at } y = 0, \quad \frac{\partial T}{\partial y} = 0$$

$$BC4: \text{ at } y = L, \quad T = T_\infty$$

## Dimensionless form

$$\theta = \frac{T - T_\infty}{(q''' L^2 / k)}, \quad X = \frac{x}{L}, \quad Y = \frac{y}{L}$$

Non-dimensionalizing using the dimensionless variables:

$$\frac{\partial^2 \theta}{\partial X^2} + \frac{\partial^2 \theta}{\partial Y^2} + 1 = 0$$

$$BC1: \text{ at } X = 0, \frac{\partial \theta}{\partial X} = 0$$

$$BC2: \text{ at } X = 1, \theta = 0$$

$$BC3: \text{ at } Y = 0, \frac{\partial \theta}{\partial Y} = 0$$

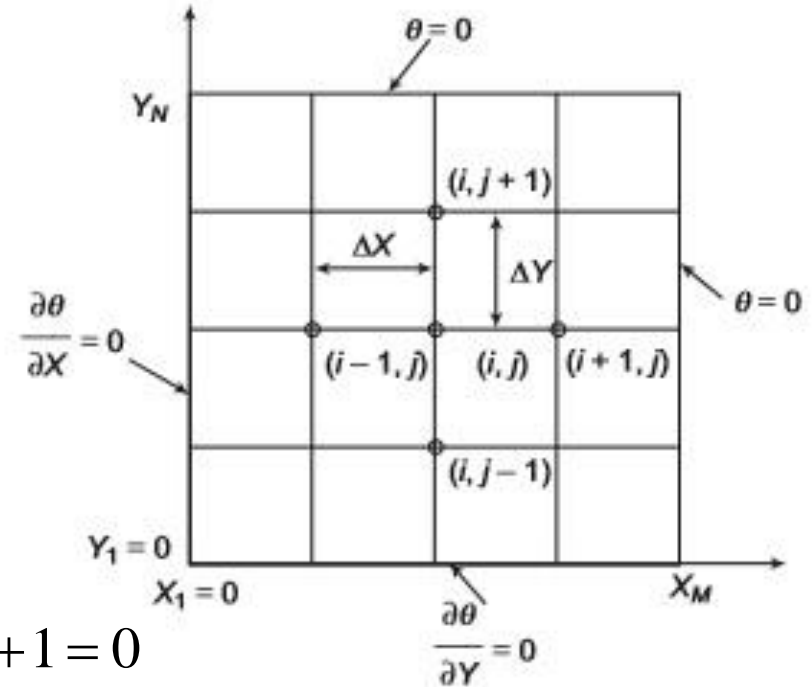
$$BC4: \text{ at } Y = 1, \theta = 0$$

## Discretization

The equation is discretized at any interior grid point  $(i, j)$  using central difference as follows:

$$\frac{\partial^2 \theta}{\partial X^2} + \frac{\partial^2 \theta}{\partial Y^2} + 1 = 0$$

$$\Rightarrow \frac{\theta_{i+1,j} - 2\theta_{i,j} + \theta_{i-1,j}}{(\Delta X)^2} + \frac{\theta_{i,j+1} - 2\theta_{i,j} + \theta_{i,j-1}}{(\Delta Y)^2} + 1 = 0$$



For a uniform grid,  $\Delta X = \Delta Y$

$$\Rightarrow -\theta_{i-1,j} - \theta_{i,j-1} + 4\theta_{i,j} - \theta_{i,j+1} - \theta_{i+1,j} = (\Delta X)^2 \quad (a)$$

❖ **Boundary condition along X = 0:**

Using image point technique  $\theta_{i-1,j} = \theta_{i+1,j}$

Setting  $i=1$  in Eq. (a), we have

$$-2\theta_{2,j} - \theta_{1,j-1} + 4\theta_{1,j} - \theta_{1,j+1} = (\Delta X)^2$$

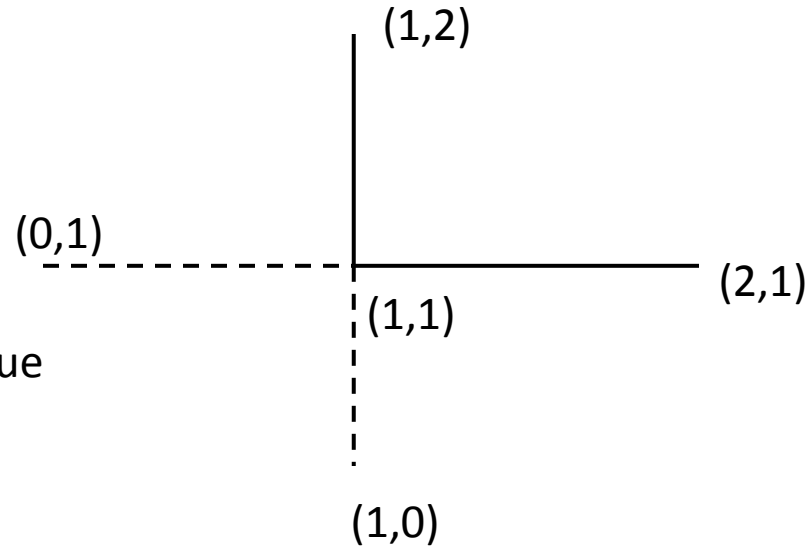
❖ **Boundary condition along Y= 0:**

Using image point technique  $\theta_{i,j-1} = \theta_{i,j+1}$

Setting  $j=1$  in Eq. (a), we have

$$-\theta_{i+1,1} - 2\theta_{i,2} + 4\theta_{i,1} - \theta_{i-1,1} = (\Delta X)^2$$

## Handling of corner points



Using image point technique

$$\theta_{0,1} = \theta_{2,1}$$

$$\theta_{1,0} = \theta_{1,2}$$

$$\left( \frac{\partial^2 \theta}{\partial X^2} \right)_{1,1} + \left( \frac{\partial^2 \theta}{\partial Y^2} \right)_{1,1} + 1 = 0$$

$$\Rightarrow \frac{\theta_{2,1} - 2\theta_{1,1} + \theta_{0,1}}{(\Delta X)^2} + \frac{\theta_{1,2} - 2\theta_{1,1} + \theta_{1,0}}{(\Delta Y)^2} + 1 = 0$$

$$\Rightarrow \frac{\theta_{2,1} - 2\theta_{1,1} + \theta_{2,1}}{(\Delta X)^2} + \frac{\theta_{1,2} - 2\theta_{1,1} + \theta_{1,2}}{(\Delta Y)^2} + 1 = 0$$

$$2\theta_{2,1} - 4\theta_{1,1} + 2\theta_{1,2} + (\Delta X)^2 = 0$$

For a uniform grid,  $\Delta X = \Delta Y$

# Solution method Gauss-Seidel method

Equations:

$$-\theta_{i-1,j} - \theta_{i,j-1} + 4\theta_{i,j} - \theta_{i,j+1} - \theta_{i+1,j} = (\Delta X)^2$$
$$-2\theta_{2,j} - \theta_{1,j-1} + 4\theta_{1,j} - \theta_{1,j+1} = (\Delta X)^2$$
$$-\theta_{i+1,1} - 2\theta_{i,2} + 4\theta_{i,1} - \theta_{i-1,1} = (\Delta X)^2$$
$$2\theta_{2,1} - 4\theta_{1,1} + 2\theta_{1,2} + (\Delta X)^2 = 0$$

Pseudo code:

*for*  $i = 1, j = 1$

$$\theta_{1,1} = \frac{1}{4} \left( 2\theta_{2,1} + 2\theta_{1,2} + (\Delta X)^2 \right)$$

*for*  $i = M, j = 2, N - 1$

$$\theta_{1,j} = \frac{1}{4} \left[ (\Delta X)^2 + 2\theta_{2,j} + \theta_{1,j-1} + \theta_{1,j+1} \right]$$

*for*  $i = 2, M - 1, j = 1$

$$\theta_{i,1} = \frac{1}{4} \left[ (\Delta X)^2 + \theta_{i+1,1} + 2\theta_{i,2} + \theta_{i-1,1} \right]$$

*for*  $i = 2, M - 1, j = 2, N - 1$

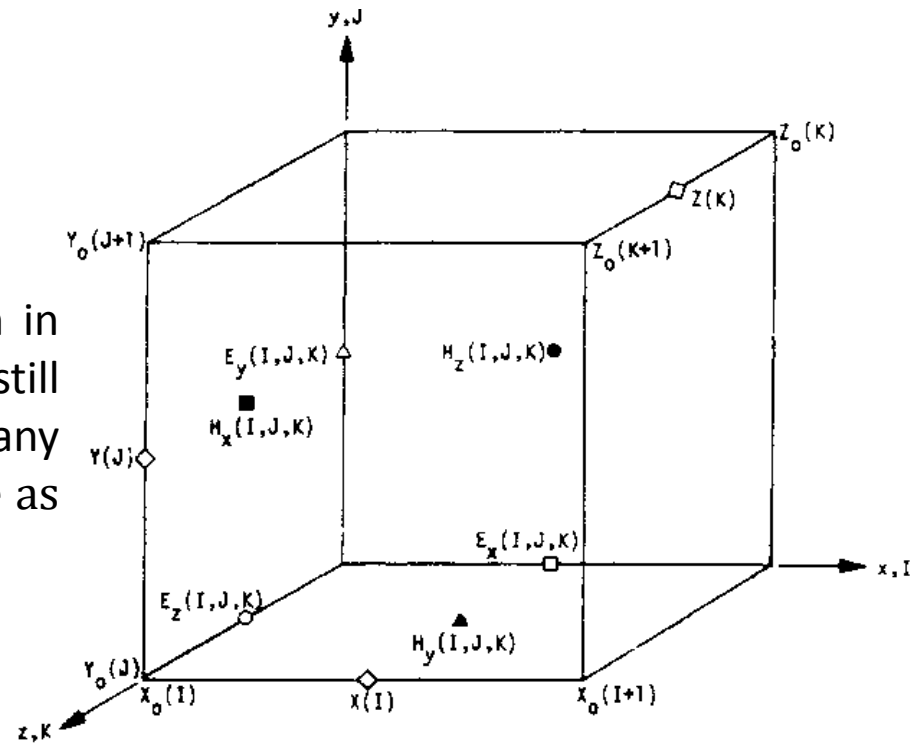
$$\theta_{i,j} = \frac{1}{4} \left[ (\Delta X)^2 + \theta_{i-1,j} + \theta_{i,j-1} + \theta_{i,j+1} + \theta_{i+1,j} \right]$$

Initial guess,  
*for*  $i = 1, M, J = 1, N$   
 $\theta_{i,j} = 0$



# Three-dimensional problems

For three-dimensional steady heat conduction in Cartesian coordinates, the basic approach is still the same. The equation is discretized at any interior grid point  $(i,j,k)$  using central difference as follows:



$$\frac{\partial^2 \theta}{\partial X^2} + \frac{\partial^2 \theta}{\partial Y^2} + \frac{\partial^2 \theta}{\partial Z^2} = 0$$

$$\Rightarrow \frac{\theta_{i+1,j,k} - 2\theta_{i,j,k} + \theta_{i-1,j,k}}{(\Delta X)^2} + \frac{\theta_{i,j+1,k} - 2\theta_{i,j,k} + \theta_{i,j-1,k}}{(\Delta Y)^2} + \frac{\theta_{i,j,k+1} - 2\theta_{i,j,k} + \theta_{i,j,k-1}}{(\Delta Z)^2} + 1 = 0$$

For a uniform grid,  $\Delta X = \Delta Y = \Delta Z$

$$\Rightarrow \theta_{i+1,j,k} + \theta_{i-1,j,k} + \theta_{i,j+1,k} + \theta_{i,j-1,k} + \theta_{i,j,k+1} + \theta_{i,j,k-1} - 6\theta_{i,j,k} + (\Delta X)^2 = 0$$